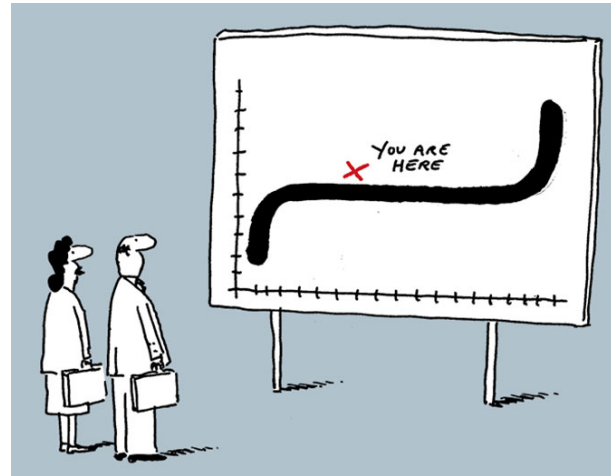


Strategies for Nuclear Resonant Experiments



Wolfgang Sturhahn

Good strategy matters:

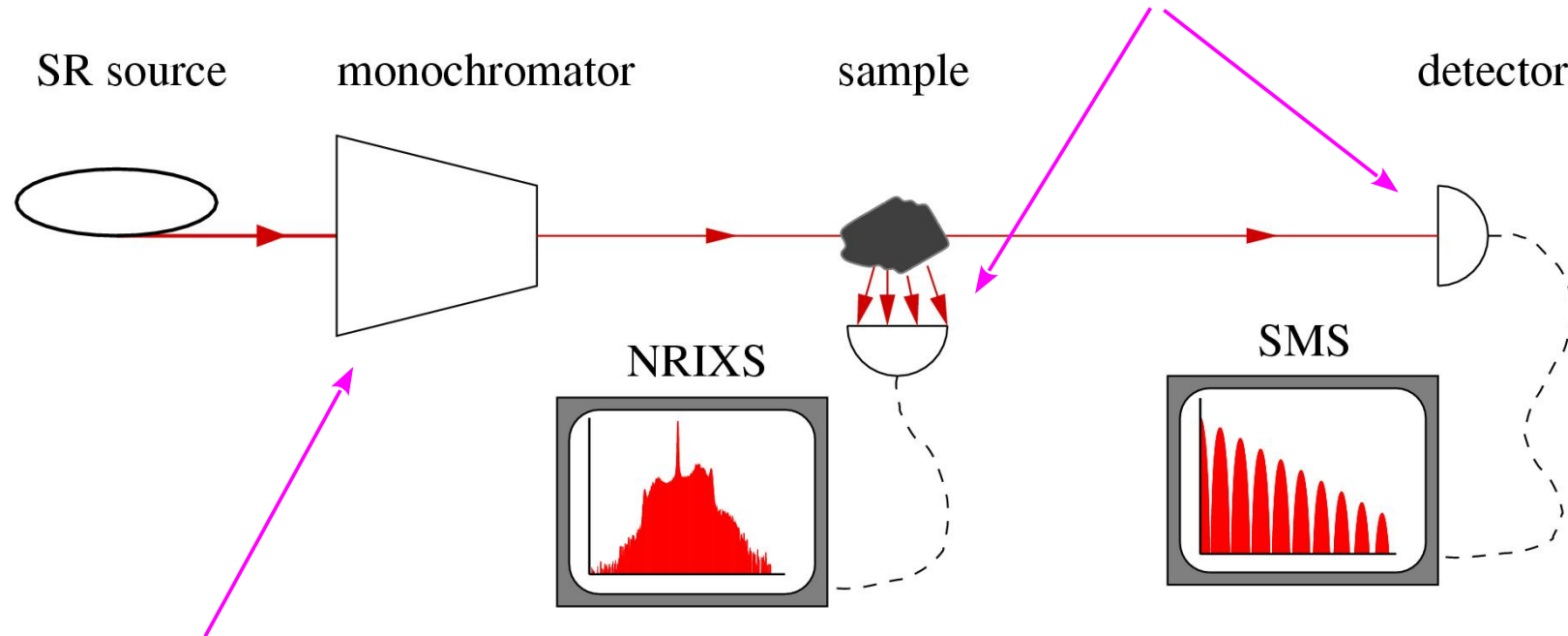
- achieve goals by following a plan over an extended time period
 - ☆ What are the scientific goals?
Can they be achieved by NRS experiments?
⇒ knowledge of science area and NRS
 - ☆ Which NRS method is appropriate for scientific goals?
Can a reasonable experimental plan be developed?
⇒ forward modeling tools
⇒ beamline capabilities
⇒ experiment readiness definition
 - ☆ How much time is needed for relevant experiments?
Are non-NRS experiments needed beforehand?
⇒ counting rate estimates, staffing requirements
 - ☆ We have the data! - Now what?
⇒ data evaluation tools



NRS, experimental setup:

- x-ray pulses must be sufficiently separated in time

- detectors must have good time resolution and excellent dynamic range



- monochromatization to meV-level required
- energy is tuned around nuclear transition (NRXS)

useful reviews of Nuclear Resonant Spectroscopy:

E. Gerdau and H. deWaard, eds., Hyperfine Interact. 123-125 (1999-2000)

W. Sturhahn, J. Phys.: Condens. Matt. 16 (2004)

R. Röhlberger, Nuclear Condensed Matter Physics with Synchrotron Radiation: Basic Principles, Methodology and Applications, Springer (2004)

W. Sturhahn and J.M. Jackson, GSA special paper 421 (2007)



The nucleus as a probe:

➤ The nucleus is not at rest:

NRIXS – Nuclear Resonant Inelastic X-ray Scattering (a.k.a. NRVS and NIS)

- ☆ energy/momentum conservation ⇒ recoil energy shift
- ☆ velocity in gases ⇒ Doppler shift
- ☆ vibrations in solids ⇒ phonon excitation/annihilation, recoilless absorption
- ☆ local vibrational density of states
- ☆ applications include determination of sound velocities and thermodynamic properties

➤ The nucleus is not a point charge:

SMS – Synchrotron Mössbauer Spectroscopy (a.k.a. NFS)

- ☆ internal dynamics ⇒ nuclear transitions
- ☆ volume ⇒ isomer shift
- ☆ spin ⇒ magnetic level splitting
- ☆ quadrupole moment ⇒ quadrupole splitting
- ☆ internal magnetic fields, electric field gradients, isomer shifts
- ☆ applications include magnetic phase transitions, determination of spin & valence states, and melting studies



What the beamline provides:

➤ Time

➤ Instruments

- ☆ monochromator with appropriate bandwidth & flux
- ☆ rotation & translation stages with appropriate range & resolution
- ☆ focusing optics for small samples
- ☆ detectors & miscellaneous electronics

➤ Software tools

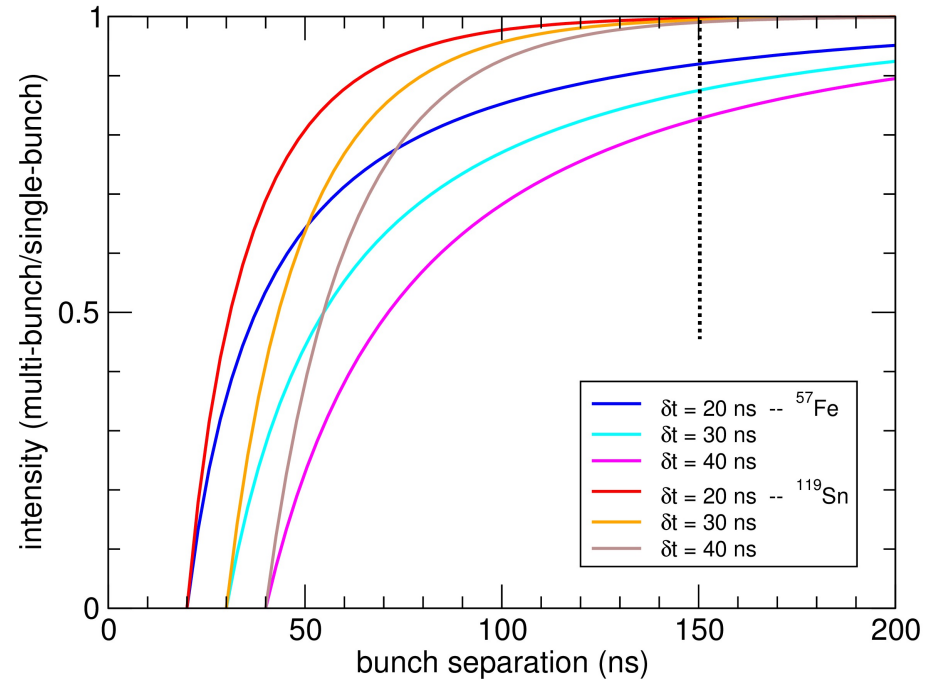
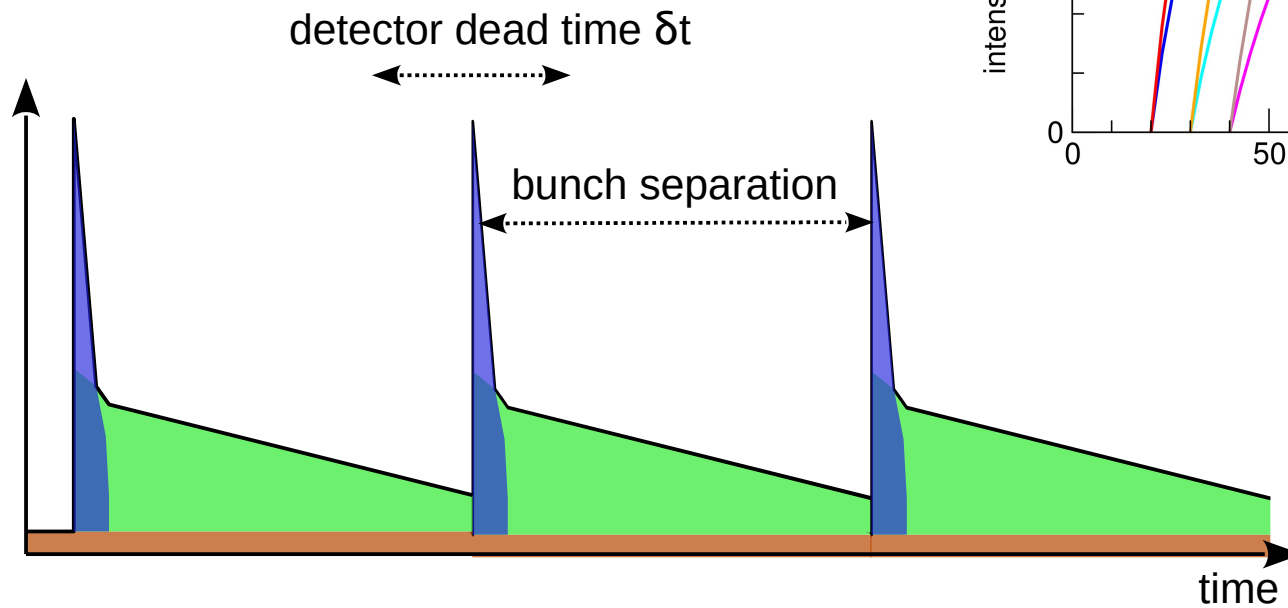
- ☆ data acquisition & motion control
- ☆ *in-situ* data evaluation for experimental progress assessment

➤ Storage ring time structure



Time structure of synchrotron radiation:

detector dead time and bunch separation determine the detectable counts in NRIXS.

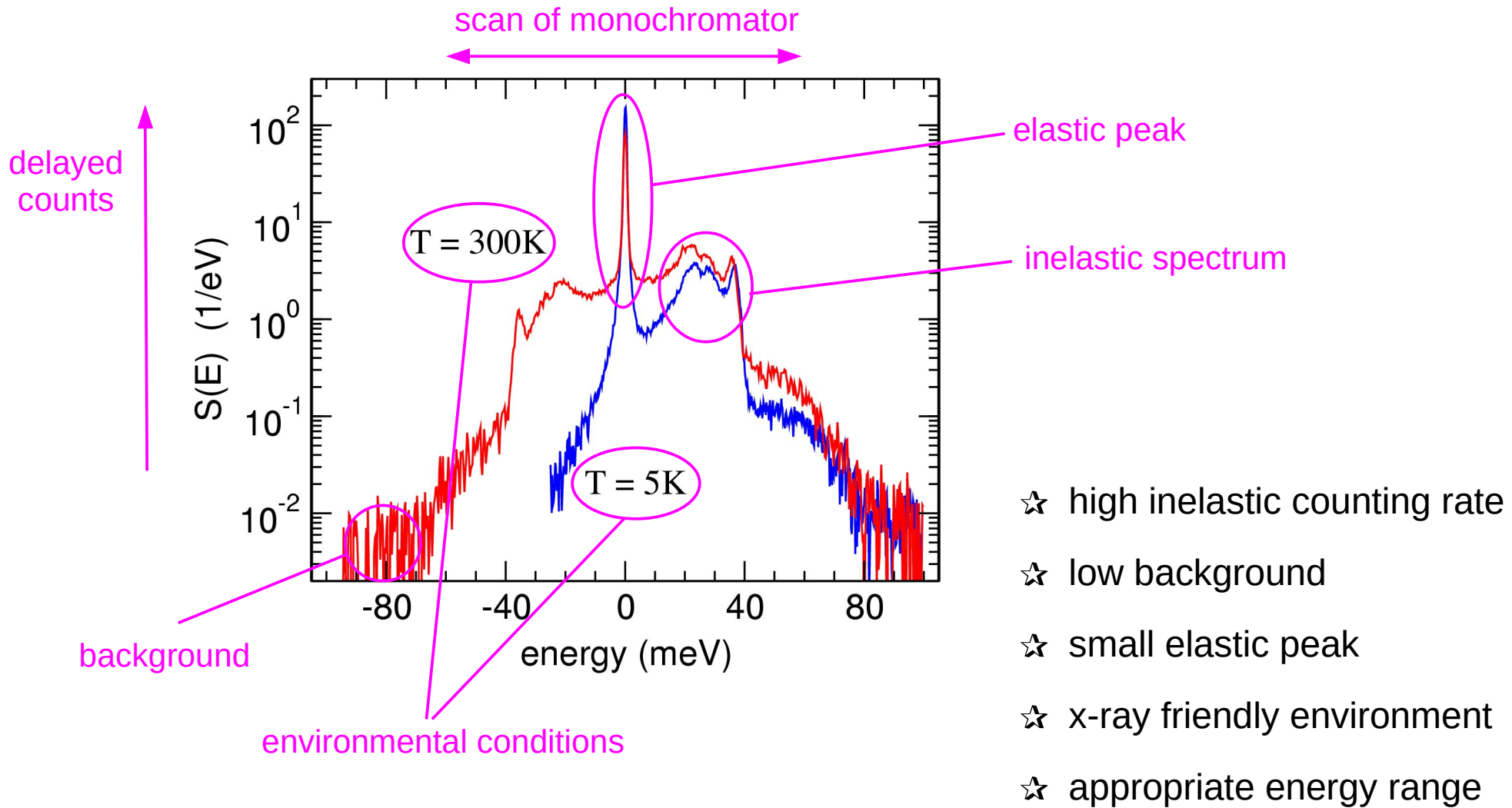


$$\tau = \hbar / \Gamma$$

141 ns for ^{57}Fe
25.6 ns for ^{119}Sn



NRIXS sample:

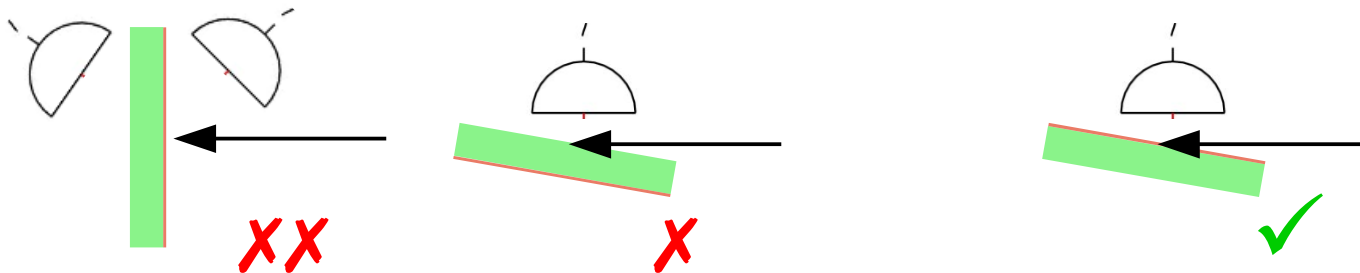


NRIXS, inelastic counting rate:

- high density of resonant isotope:
enrichment often necessary, avoid short lifetimes

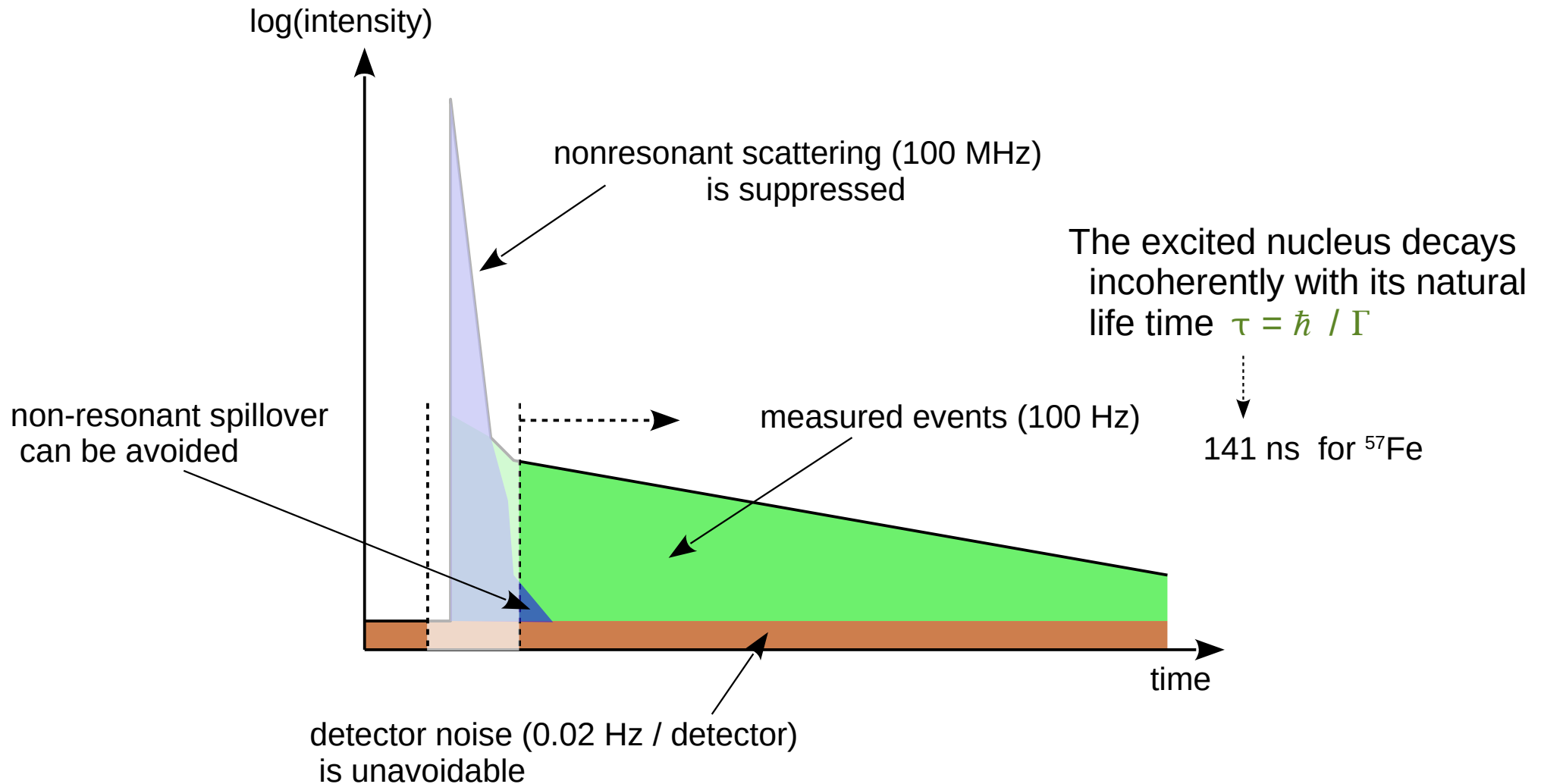
☆ ^{57}Fe :	2.2 %	141 ns
☆ ^{83}Kr :	11.6 %	212 ns
☆ ^{119}Sn :	8.6 %	25.6 ns
☆ ^{151}Eu :	47.8 %	13.8 ns
☆ ^{161}Dy :	18.9 %	40.7 ns

- scattering geometry: large collection angle, large scattering volume



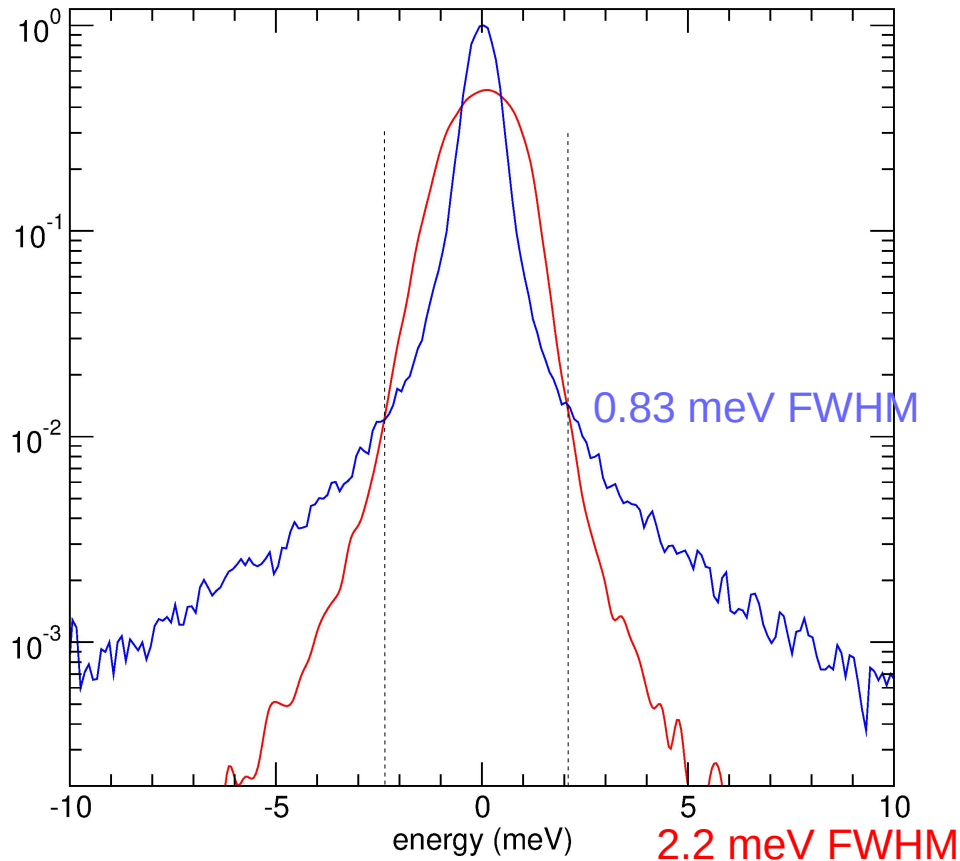
NRIXS, low background:

- background has two contributions:
detector noise, non-resonant spillover



NRIXS, elastic peak:

- choice of monochromator: bandwidth, shape, flux, stability

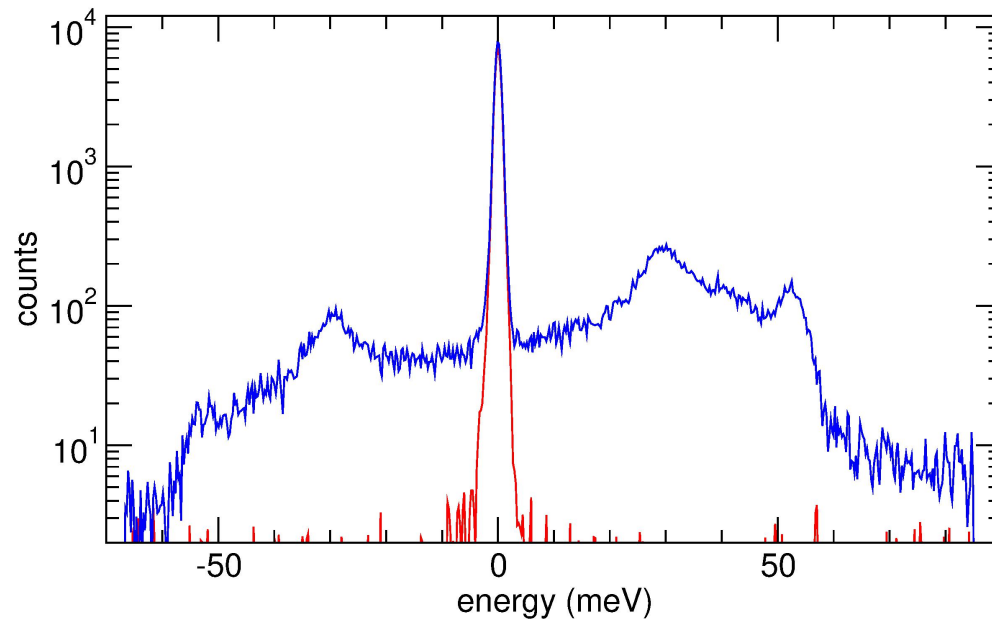
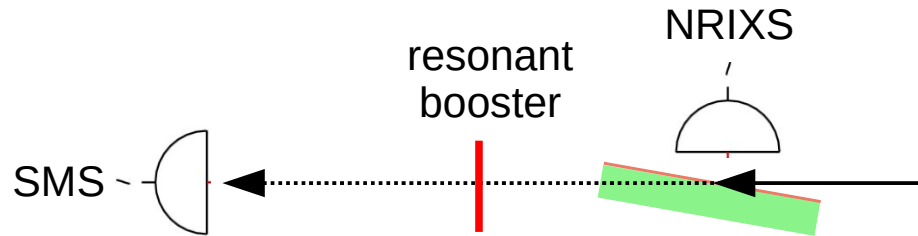



- ★ monochromator design influences bandwidth and tails in different ways
- ★ smaller tails are very important for weak scatterers such as:
 - thin samples
 - dilute samples
 - low isotopic abundance



NRIXS, elastic peak:

- mitigate stability issues by *in-situ* measurement of mono response

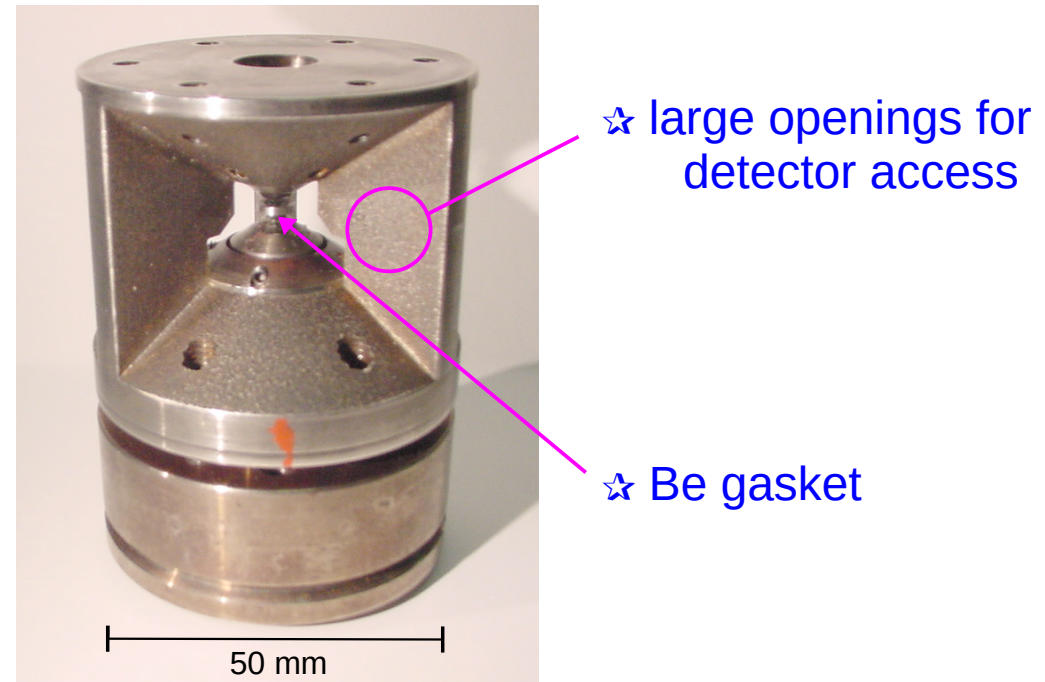
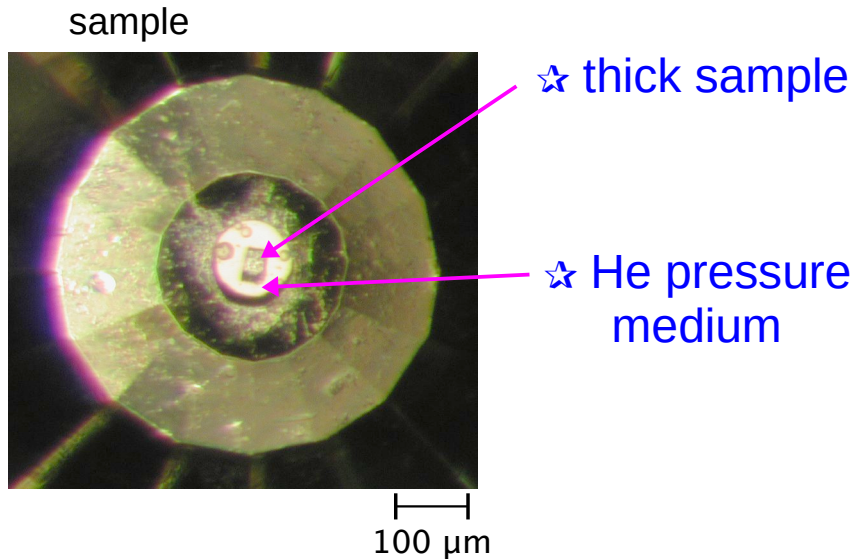


- ☆ large distance between sample and SMS detector
- ☆ SMS signal \geq NRIXS peak-signal
- ☆ use resonant booster only if sample is homogeneous  over x-ray profile

NRIXS, environment:

- sample environment has to be “x-ray friendly”:
low absorption for observed x-rays

☆ ^{57}Fe :	14.4 keV	6.4 keV
☆ ^{83}Kr :	9.3 keV	1.6 keV
☆ ^{119}Sn :	23.9 keV	3.4 keV
☆ ^{151}Eu :	21.5 keV	5.8 keV
☆ ^{161}Dy :	25.7 keV	6.5 keV



NRIXS, energy range:

- detailed balance renders NRIXS spectrum asymmetric
 - ☆ start scan between -100 meV and -80 meV at room temperature
 - ☆ end scan depends on sample structure, pressure, temperature
- exclusion of negative energies completely disables checks on quality and consistency of data **XXX**
- recommended end scan energies:
 - ☆ metals and alloys 80 - 120 meV
 - ☆ oxides and silicates without hydrogen 100 - 150 meV
 - ☆ oxides and silicates with hydrogen 150 - 200 meV
- selecting too small end scan energies excludes high-energy vibrations and compromises data quality **XXX**



Interpretation of NRIXS spectra:

- NRIXS spectra directly provide the Fourier transform of the self-intermediate scattering function

$$S(\mathbf{k}, E) = \frac{1}{2\pi\hbar} \int \left\langle e^{i\mathbf{k}\hat{\mathbf{r}}(t)} e^{-i\mathbf{k}\hat{\mathbf{r}}(0)} \right\rangle e^{iEt/\hbar} dt$$

- In the quasi-harmonic approximation the partial projected phonon density-of-states is obtained by a multi-phonon expansion

$$S(\mathbf{k}, E) = f(\mathbf{k})\delta(E) + \sum_{n=1}^{\infty} S_n(\mathbf{k}, E)$$
$$S_1(\mathbf{k}, E) = f(\mathbf{k}) \frac{E_R}{E(1 - \exp[-\beta E])} g(\mathbf{k}, |E|)$$
$$S_n(\mathbf{k}, E) = \frac{1}{n f(\mathbf{k})} \int S_{n-1}(\mathbf{k}, E') S_1(\mathbf{k}, E - E') dE'$$
$$f(\mathbf{k}) = \exp \left[- \int \frac{E_R}{E} \coth\left(\frac{\beta E}{2}\right) g(\mathbf{k}, E) dE \right]$$

W.Sturhahn and V.G.Kohn, *Hyperfine Interact.* 123/124 (1999)



Information from NRIXS spectra:

➤ directly from the data, $S(E)$

⇒ temperature

$$T = -\frac{E}{k_B} \ln \left[\frac{S(-E)}{S(E)} \right]$$

⇒ mean square displacement

$$\langle u^2 \rangle = -\frac{1}{k^2} \ln \left[1 - \int \{S(E) - S(0)\} dE \right]$$

⇒ kinetic energy

$$E_{kin} = \frac{1}{4E_R} \int (E - E_R)^2 S(E) dE$$

⇒ average force constant

$$D = \frac{k^2}{2E_R^2} \int (E - E_R)^3 S(E) dE$$

k ~ wave number of nuclear transition

E_R ~ recoil energy

ρ ~ mass density

➤ quasi-harmonic lattice model

⇒ partial phonon density of states

$$\mathcal{D}(E)$$

⇒ Debye sound velocity

$$v_D = \left(\frac{M}{2\rho\pi^2\hbar^3} \frac{E^2}{\mathcal{D}(E \rightarrow 0)} \right)^{1/3}$$

⇒ Grüneisen parameter

$$\gamma_D = \frac{1}{3} + \frac{\rho}{v_D} \left(\frac{\partial v_D}{\partial \rho} \right)_T$$

⇒ isotope fractionation

$$\ln \beta = -\frac{\Delta m}{M} \frac{1}{8(k_B T)^2} \int E^2 \mathcal{D}(E) dE$$

M ~ mass of resonant isotope

Δm ~ isotope mass difference

k_B ~ Boltzmann's constant

T ~ temperature



NRIXS evaluation, the PHOENIX GUI:

- GUI upgrade, PHOENIX-3.x, supported by Caltech
 - ☆ translates functionality into Tcl/Tk for Unix and MacOS
 - ☆ maintains all previous capabilities of CLI
 - ☆ enhancements of core modules
 - ☆ cross-project analysis tools

- upgrades to core modules
 - ☆ consistency optimization
 - ☆ advanced elastic peak subtraction
 - ☆ probability distribution analysis for Debye velocity determination

- Thanks to Jennifer Jackson and her group at Caltech for continuous tests of the software and for ongoing discussions for improvements

The screenshot displays the PHOENIX GUI interface. The left panel shows 'Phox input parameters' for a 57Fe sample, including file names, nuclear isotope, and various fitting and background options. The right panel shows 'Phox runtime messages' with a list of derived quantities and consistency tests. The bottom right shows the 'Project directory' containing files like Data/, Fe_cvm.bin, and Fe_sum.dat.

```
peak : 1.00E+00
valid range (meV) : -9.90 to 9.90
fit range (meV) : -2.99 to 2.99 centered at 0.11
ChiA2 peak FWHM pos. %bgr. asm.
start: 354.90 9.03E+03 1.20 0.107 3.107 1.08
final: 2.94 9.50E+03 1.00 0.013 1.833 1.08
errors: 0.46 3.21E+02 0.04 0.031 0.863 0.11

smoothing range (meV) : -2.897 to 3.084

Quantities derived directly from the data -->
Lamb-Moessbauer factor : 0.7995 +- 0.0019
kinetic energy : 14.0707 +- 0.0823 meV/atom
quantum excess energy : 1.2739 +- 0.0823 meV/atom
mean force constant : 175.5276 +- 3.7350 N/m

Quantities derived after refinement -->
Lamb-Moessbauer factor : 0.7996 +- 0.0019
kinetic energy : 14.0804 +- 0.0824 meV/atom
quantum excess energy : 1.2837 +- 0.0824 meV/atom
mean force constant : 177.1053 +- 3.7686 N/m
isotope fractionation : 1.5047 +- 0.0986 perMille/%
high T isotope frac. : 1.5469 +- 0.0329 perMille/%

Consistency tests using the data -->
tested quantity %deviation norm.dev. status
detailed balance 0.05 +- 0.53 0.09 ok
energy/temp. calib. -1.15 +- 0.86 1.34 acceptable

Consistency tests using the phonon DOS -->
tested quantity %deviation norm.dev. status
norm of DOS 0.04 +- 0.43 0.10 ok
Lamb-Moessbauer factor 0.01 +- 0.26 0.02 ok
kinetic energy per atom 0.13 +- 0.75 0.17 ok
mean force constant 0.81 +- 2.51 0.32 ok

deviator = quadr. mean of norm.dev. --> 0.68 ok

Quantities calculated from the partial DOS -->
Lamb-Moessbauer factor : 0.7997 +- 0.0008
kinetic energy : 14.0982 +- 0.0656 meV/atom
mean force constant : 178.5330 +- 2.3673 N/m
isotope fractionation : 1.5255 +- 0.0769 perMille/%
high T isotope frac. : 1.5594 +- 0.0207 perMille/%
Lamb-Moessbauer factor at T=0 : 0.9235 +- 0.0003
kinetic energy at T=0 : 6.8901 +- 0.0442 meV/atom
vibrational specific heat : 2.7103 +- 0.0111 k_B/atom
vibrational entropy : 3.0726 +- 0.0118 k_B/atom
resilience : 104.3601 +- 0.5060 N/m
Lamb-Moessbauer temperature : 1416.9 +- 6.87 K

-- CPU time : user 0.77 s system 0.01 s
-- PHOENIX module PHOX Finished
```


SMS sample:

- SMS is a transmission spectroscopy
 - ☆ What is the optimal sample thickness?
 - ☆ What counting rate can be expected?

- use CONUSS software to model sample signal
 - ☆ sample composition and density
 - ☆ expected hyperfine parameters
 - ☆ useful time range of measured spectrum

- use CONUSS software to evaluate measured data



CONUSS software provides solutions:

problem	program	SIF	examples
fitting data forward scattering dual fit Mössbauer spectroscopy grazing incidence Bragg/Laue diffraction	kctl	in_kctl in_kfor in_kfor in_kfor in_kgin in_kref	kctl-NFS1, kctl-NFS2 kctl-NFS3 kctl-MBS1, kctl-MBS2 kctl-GINS kctl-NBS1, kctl-NBS2
explore parameter space forward scattering or Mössbauer grazing incidence Bragg/Laue diffraction	kmco	in_kmco in_kfor in_kgin in_kref	kmco-NFS kmco-GINS kmco-NBS
calculate spectra forward scattering or Mössbauer grazing incidence Bragg/Laue diffraction	kfmf kgmf krmf	in_kfor in_kgin in_kref	kfmf-NFS, kfor-NFS kgmf-GINS, kgmf-GIS krmf-NBS

Interpretation of SMS spectra:

- Nuclear resonant contribution to the index-of-refraction

$$\delta \mathbf{n}(E) = \frac{\Gamma}{4k} F_{LM} \sigma_{0\rho} \sum_{mm'} \frac{\mathbf{W}_{mm'}}{E_{mm'} - E - i\Gamma/2}$$

- Time spectrum

$$\frac{dI}{dt} = \left| \int \left[e^{ikD\delta \mathbf{n}(E)} - 1 \right] e^{-iEt/h} \frac{dE}{2h} \right|^2$$

- Mössbauer transmission spectrum

$$T(E) = \int \text{Trace} \left[e^{-kD\text{Im}[\delta \mathbf{n}(E')]} \right] L(E - E') dE'$$

W.Sturhahn, J.Phys.: Condens.Matt. 16 (2004)



SMS sample thickness:

- ☆ single line:
 - isomer shift only
- ☆ two lines:
 - electric field gradient, quadrupole splitting
 - two sites with different isomer shifts
- ☆ many lines:
 - magnetic field
 - several sites with different line positions

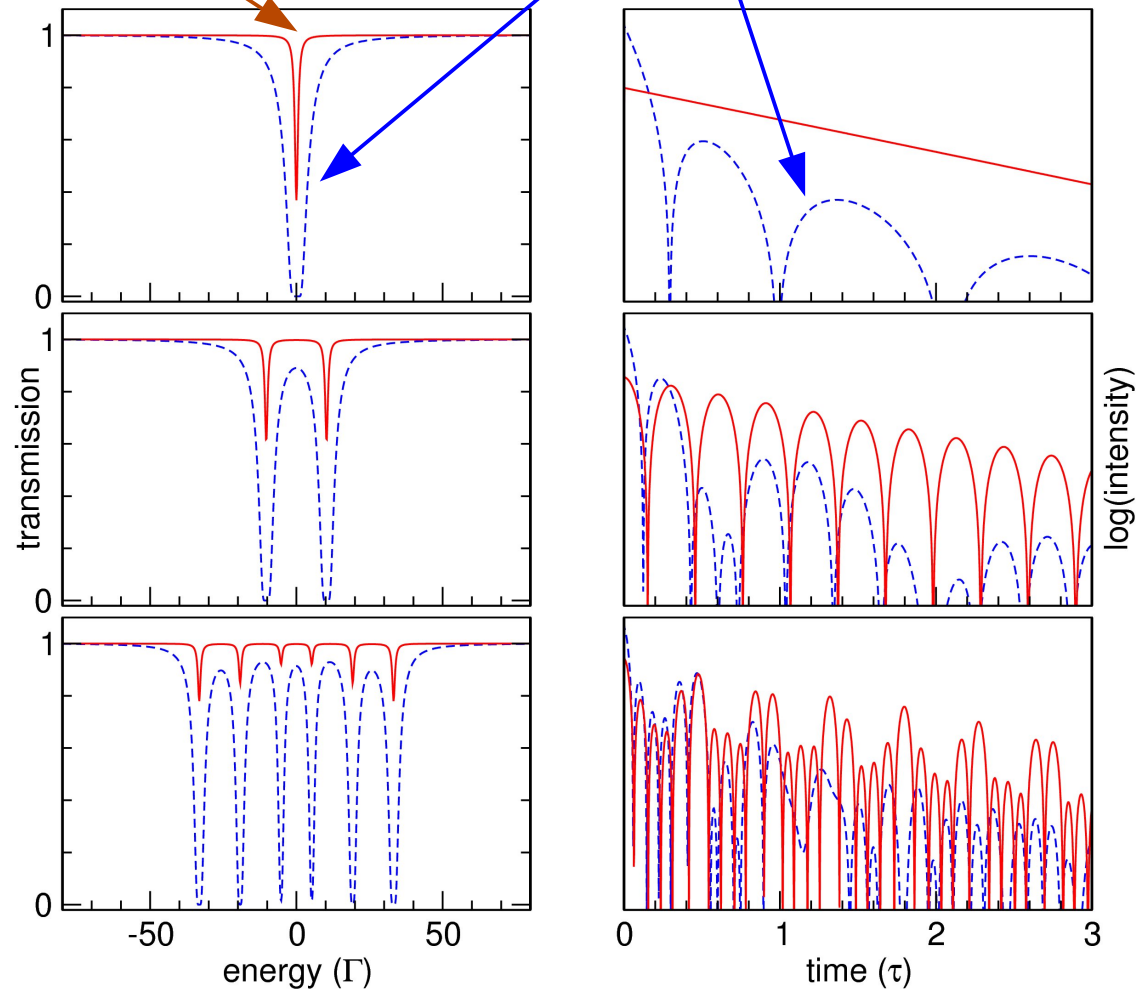
effective thickness:

$$D_{\text{eff}} = F_{\text{LM}} \sigma_0 \rho D$$

Lamb-Mössbauer factor
 resonant cross section
 nuclei per area
 geometric thickness

undisturbed line shape, $D_{\text{eff}} = 1$

line broadening, $D_{\text{eff}} = 50$



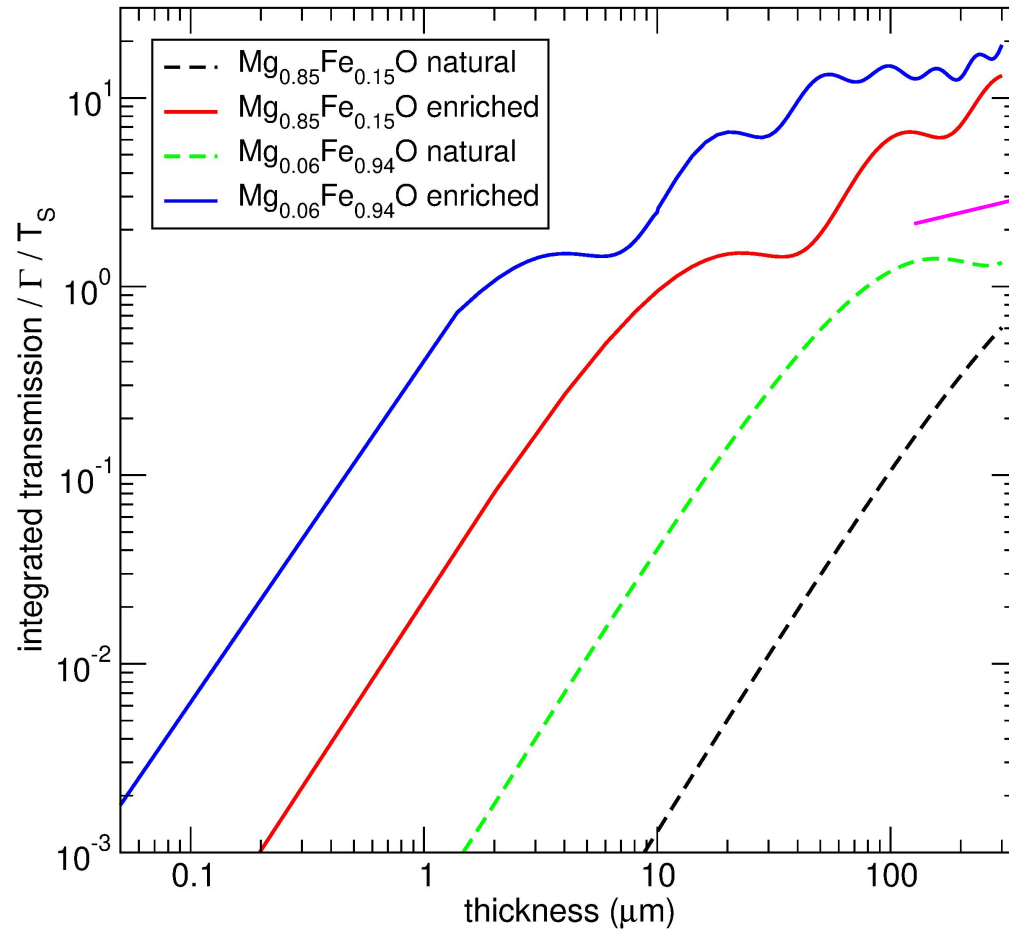
Mössbauer spectroscopy

SMS

☆ calculations performed with CONUSS software



SMS integrated transmission:



☆ calculations performed with CONUSS software

➤ delayed flux

$$F = \frac{N}{\Gamma T_S} T_S T_A \epsilon_D I_\Gamma$$

➤ prompt flux may saturate detector

$$P = T_S T_A \epsilon_D I_0 \leq P_{sat}$$

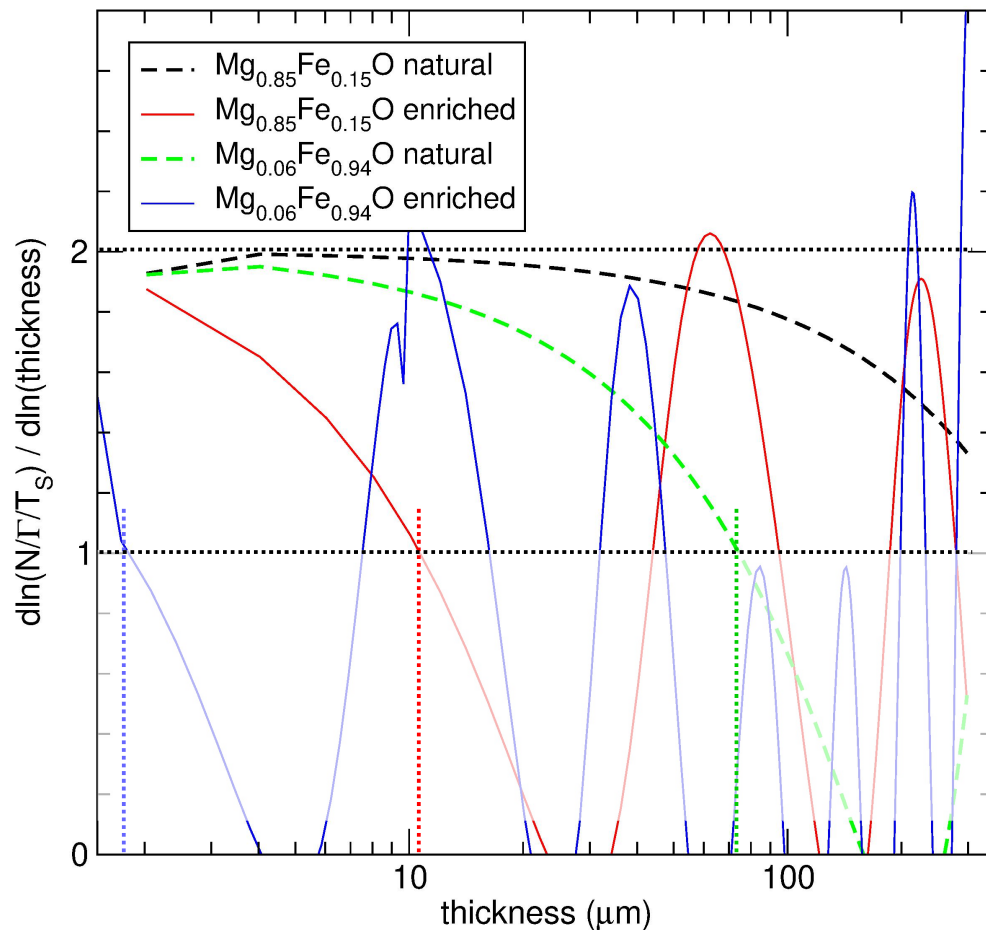
➤ maximum delayed flux

$$F \leq \frac{N}{\Gamma T_S} \frac{\Gamma}{\Delta E} P_{sat}$$

≈ 500 for ⁵⁷Fe



SMS choice of thickness:



- the logarithmic derivative

$$\frac{d}{d \ln D} \frac{N}{\Gamma T_S}$$

decreases with thickness
from values close to 2.

- values > 1 usually mean acceptable thickness effects
- maximum counting rates

$\text{Fe}_{0.15}\text{Mg}_{0.85}\text{O}$	300 1/s	300 micron
$^{57}\text{Fe}_{0.15}\text{Mg}_{0.85}\text{O}$	500 1/s	11 micron
$\text{Fe}_{0.94}\text{Mg}_{0.06}\text{O}$	460 1/s	75 micron
$^{57}\text{Fe}_{0.94}\text{Mg}_{0.06}\text{O}$	450 1/s	1.5 micron



In conclusion:

- develop a good strategy by familiarizing yourself with the specifics of NRIXS and SMS
- use forward modeling tools provided with the PHOENIX and CONUSS software to obtain quantitative information for your samples
- optimized samples for NRIXS and SMS are very different
- follow the guidelines provided in this presentation
- start communications early with the staff at the beamline that you chose for your experiments
- evaluate your data with the PHOENIX and CONUSS software if possible in-situ during the experiment



Ende